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## Molecular Crystals and Liquid Crystals Science and Technology. Section A. Molecular Crystals and Liquid Crystals

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### Dipping-Induced In-Plane Molecular Alignment of LB Films

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## DIPPING-INDUCED IN-PLANE MOLECULAR ALIGNMENT OF LB FILMS

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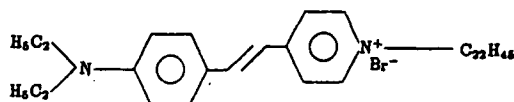
**Abstract** Super-quadratic growth of SHG intensity with thickness could be realized due to the extra dipping-induced and epitaxy-enhanced in-plane polarization in the hemicyanine LB multilayers.

### INTRODUCTION

In surface physics, one usually assumed that molecules adsorbed on substrates with an in-plane isotropy.<sup>1</sup> In fact, flow orientation during transfer of Langmuir monolayer<sup>2,3</sup> would not only modify some concepts and theories in LB technology and nonlinear optics, but also find potential applications in areas such as liquid crystal aligning agent<sup>4</sup> or waveguide frequency-doubler.<sup>5</sup>

### SAMPLE PREPARATION

The molecular structure of the hemicyanine (HEM) is:



Arachidic acid  $C_{19}H_{39}COOH$  (AA) was used as inert spacer molecules to form stable noncentrosymmetric structures. Four Y-type HEM/AA interleaving multilayer samples A, B, C, D of bilayer numbers  $N = 6, 12, 24, 36$  respectively were prepared at a surface pressure of  $30 \text{ mN m}^{-1}$  and dipping speed of  $3 \text{ mm min}^{-1}$ .

## DIPPING-INDUCED ANISOTROPY

Fig.1(a) shows the normal-incidence polarized UV-visible absorption spectra of sample D with the incident field parallel (broken line) or perpendicular (solid line) to the dipping direction on the sample. The linear anisotropy parameter was defined  $r_L \equiv A_{\parallel}/A_{\perp}$ , where  $A_{\parallel}$ ,  $A_{\perp}$  were the areas below the respective absorption curves. We obtained  $r_L = 1.66 > 1$  from Fig.1(a) revealing a significant flow orientation along the dipping direction.

The second harmonic generation (SHG) experiments were performed by using a 40ps pulsewidth, 10Hz repetition rate beam at  $1.064 \mu\text{m}$  with different angles  $\phi$  between the dipping direction and incident plane by rotating the sample around its surface normal, at an incident angle  $\theta = 45^\circ$  unless otherwise stated. The transmitted SHG signals at 532nm were detected by a photomultiplier tube and a boxcar averager. The measured p-in/p-out SHG intensity  $I_{pp}(\phi)$  from sample D is shown by the solid dots in Fig.1(b) displaying an asymmetric "oval" pattern with a maximum at  $\phi = 0^\circ$ .

Our classical nonlinear oscillator model derived<sup>6</sup>

$$I_{pp} \propto [d_1 \cos\theta \cos\phi + d_2 \sin\theta]^4 \quad (1)$$

The fit result using Eq.(1) is given by the solid line in Fig.1(b) and qualitatively agreed with the experimental data.

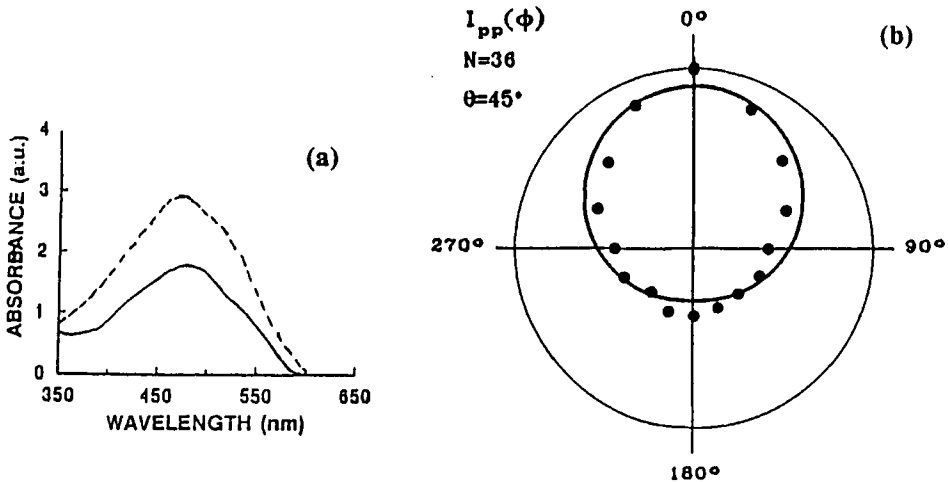


FIGURE 1 Sample D: (a) Polarized absorption spectra; (b) SHG intensity  $I_{pp}(\phi)$ .

### EPITAXY-ENHANCED ANISOTROPY

For samples A, B, C, D, we obtained  $r_L = 1.10, 1.24, 1.55, 1.66$  by absorption measurements. We also define a nonlinear anisotropy parameter  $r_{NL} \equiv [I_{pp}(\phi = 0^\circ) / I_{pp}(\phi = 180^\circ)]^{1/2}$ , and obtained  $r_{NL} = 1.07, 1.22, 1.48, 1.77$  measured at  $\Theta = 45^\circ$  and  $r_{NL} = 1.06, 1.26, 1.46, 1.68$  measured at  $\Theta = 30^\circ$  by SHG experiments. As shown in Fig.2 (a), the degree of anisotropy  $r_L$  and  $r_{NL}$  increased monotonically and consistently with  $N$ , due to the interlayer interaction or "epitaxial evolution" process during deposition:<sup>7</sup>

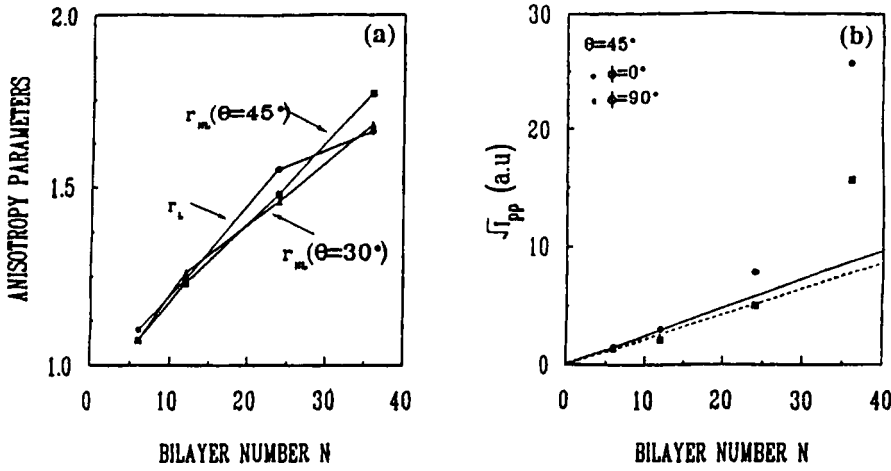


FIGURE 2 N-dependence of (a) Anisotropy parameters  $r_L$  and  $r_{NL}$ ; (b)  $\sqrt{I_{pp}}$  measured at  $\Theta = 45^\circ$ .

### SUPER-QUADRATICALLY ENHANCED SHG INTENSITY

From nonlinear optics, the SHG intensity

$$I_{2\omega} = \frac{2\omega^2 D_{\text{eff}}^2 I_{\omega}^2 L^2 \text{sinc}^2(L\Delta k/2)}{c^3 \epsilon_0 n_{\omega}^2 n_{2\omega}} \quad (2)$$

The  $L^2$ -dependence has been considered as an criterion and upper limit for perfectly and invariantly structured multilayer films within the quasi-phase-matching range of optical path  $L$  [ $L\Delta k \ll 1$ , i.e.  $\text{sinc}(L\Delta k/2) \approx 1$ ].

In the presence of the extra  $L$ -dependent polarization  $d_h$ , however, the effective second-order susceptibility  $D_{\text{eff}}$  in Eq.(2) should be replaced by

$$D_{\text{eff}}'(L) = D_{\text{eff}} [ 1 + d_h(L) \cot\theta \cos\phi / d_z ]^2 . \quad (3)$$

As  $L$  (or  $N$ ) increased, an increase of  $d_h(L)$  speeds up the intensity growth and a super-quadratic enhancement could be observed, especially for  $\phi=0^\circ$ , if the intensity reduction due to structural deterioration, low transfer ratio and phase mismatching were insignificant.

Fig.2(b) gives the data for  $\sqrt{I_{\text{pp}}}$  with  $\phi=0^\circ$  (soid dots) and  $\phi=90^\circ$  (solid squares) vs  $N$ . Super-linear growth of  $\sqrt{I_{\text{pp}}}$ , i.e. Super-quadratic increase of  $I_{\text{pp}}$  with  $N$  were clearly seen by the data points well above their respective straight lines, especially for  $\phi=0^\circ$  and larger  $N$ 's.

## ACKNOWLEDGEMENTS

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